arXiv:1301.7435 PRL 110, 166802

# Topological charge pumping of cold atoms

Lei Wang ETH Zurich

#### Collaborators:

Alexey Soluyanov Matthias Troyer

Xi Dai





## Plan

Topological charge pumping in a 1D optical lattice

## Plan

Topological charge pumping in a 1D optical lattice

Measure topological index of 2D optical lattices

# Pumps

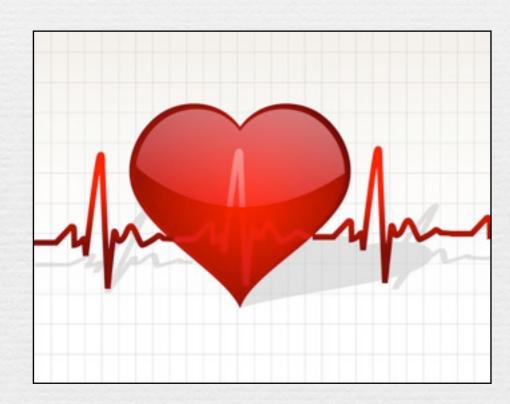


A pump is a device that moves fluids, or sometimes slurries, by mechanical action.

# Pumps



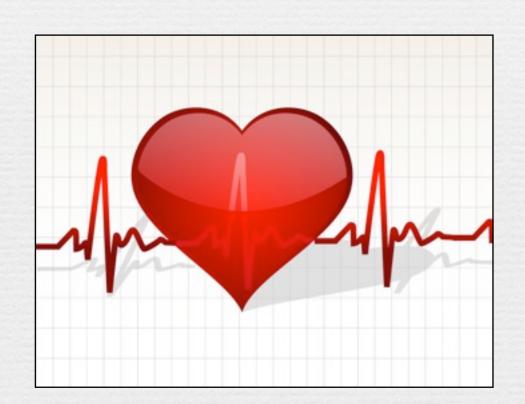
A pump is a device that moves fluids, or sometimes slurries, by mechanical action.

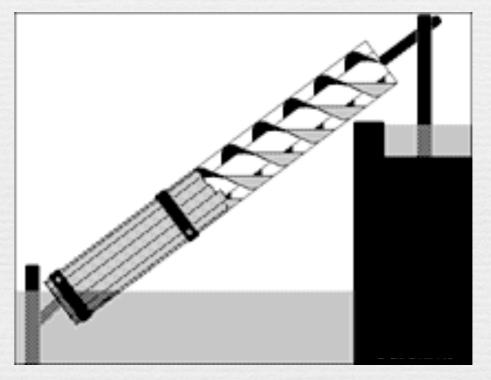


# Pumps



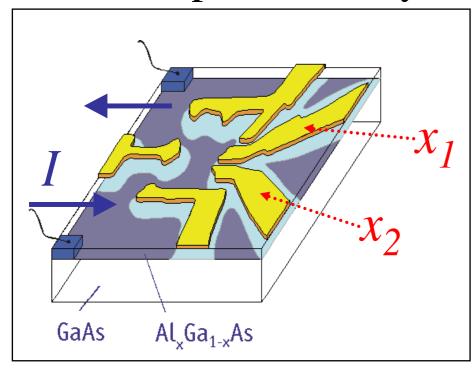
A pump is a device that moves fluids, or sometimes slurries, by mechanical action.

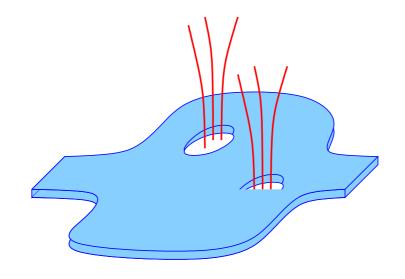




Archimedes' screw ~250 BC

 $\bullet$  conductor penetrated by Aharonov-Bohm fluxes  $\boldsymbol{x}$  ,  $\boldsymbol{x}$ 

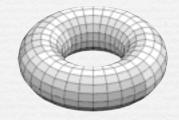




Switkes et al 1999

Archimedes' screw ~250 BC

#### Topological pump



A device transfers quantized charge in each pumping cycle.

Thouless 1983

$$H = \sum_{i} (t + \delta t) c_{Ai}^{\dagger} c_{Bi} + (t - \delta t) c_{Ai+1}^{\dagger} c_{Bi} + H.c.$$
Su, Schrieffer, Heeger, 197
$$\delta t > 0$$

$$\delta t < 0$$

$$\delta t < 0$$

$$Na_x dissipation t) + (t - \delta t) \cos ka$$

Dynamical analog of quantum Hall effect

$$d_{z}(k) = 0$$

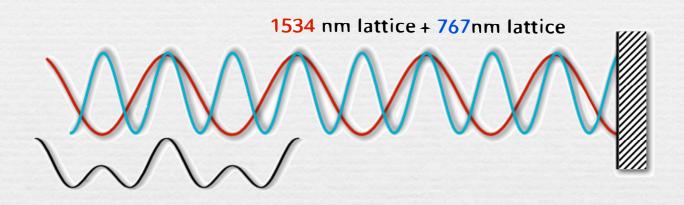
## Experimental progresses

#### Optical Superlattice

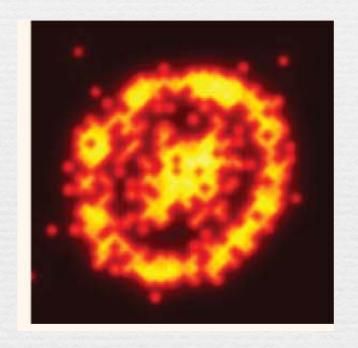
Fölling et al, Atala et al

#### in-situ imaging

Gemelke, et al, Sherson et al, Bakr et al



$$V_{\rm OL}(x) = V_1 \cos^2\left(\frac{2\pi x}{d}\right) + V_2 \cos^2\left(\frac{\pi x}{d} - \varphi\right)$$

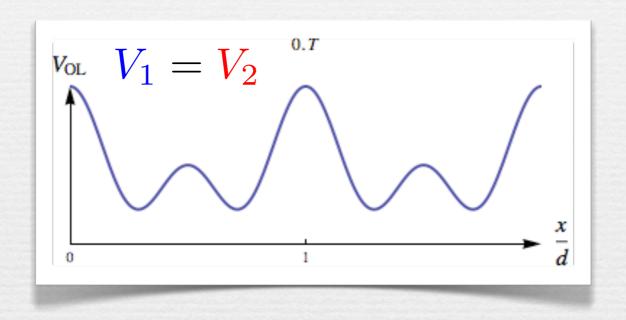


nc ca rc Allows to measure exact quantization of pumped charge

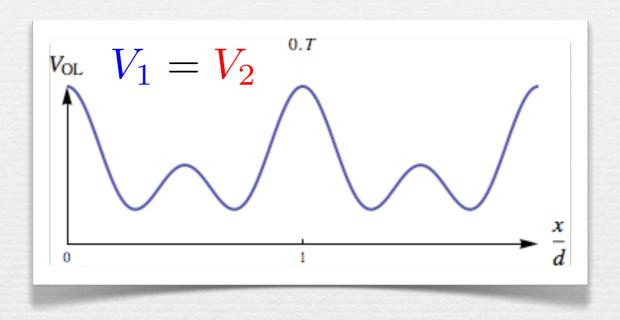
$$V_{\rm OL}(x,t) = V_1 \cos^2\left(\frac{2\pi x}{d}\right) + V_2 \cos^2\left(\frac{\pi x}{d} - \frac{\pi t}{T}\right)$$

$$V_1 = V_2$$

$$V_{\rm OL}(x,t) = V_1 \cos^2\left(\frac{2\pi x}{d}\right) + V_2 \cos^2\left(\frac{\pi x}{d} - \frac{\pi t}{T}\right)$$



$$V_{\rm OL}(x,t) = V_1 \cos^2\left(\frac{2\pi x}{d}\right) + V_2 \cos^2\left(\frac{\pi x}{d} - \frac{\pi t}{T}\right)$$

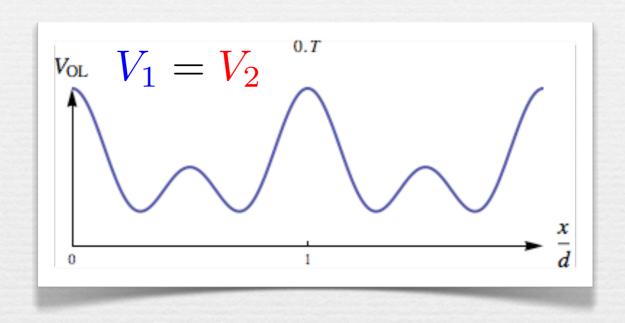


0 A = B - A = B

Su, Schrieffer, Heeger, 1979

$$T/2$$
 A — B = A — B

$$V_{\rm OL}(x,t) = V_1 \cos^2\left(\frac{2\pi x}{d}\right) + V_2 \cos^2\left(\frac{\pi x}{d} - \frac{\pi t}{T}\right)$$



$$0 A = B - A = B$$

T/4 A ---- B ---- A ---- B

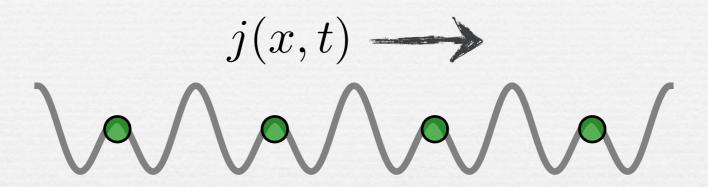
T/2 A — B = A — B

3T/4 A ---- B ---- B

Su, Schrieffer, Heeger, 1979

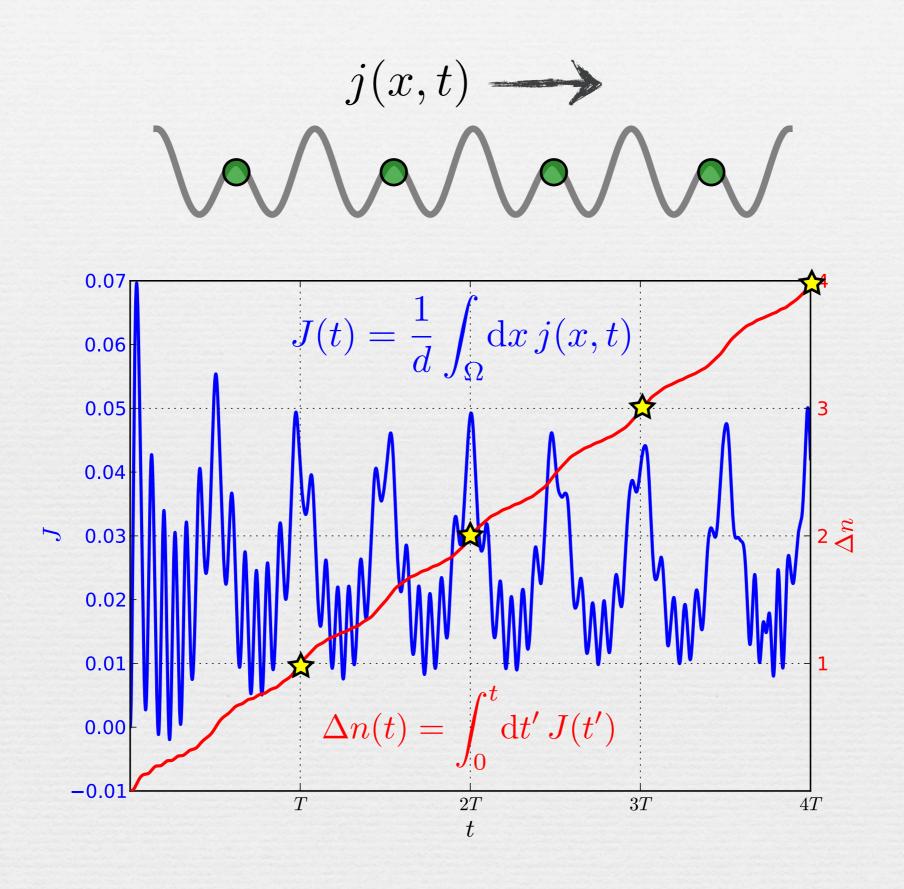
Rice, Mele, 1982

#### Pumping dynamics



$$H(x,t) = -\frac{\hbar^2}{2m} \nabla^2 + V_{\rm OL}(x,t)$$
$$i\frac{\partial}{\partial t} |\Psi\rangle = H(x,t) |\Psi\rangle$$

#### Pumping dynamics

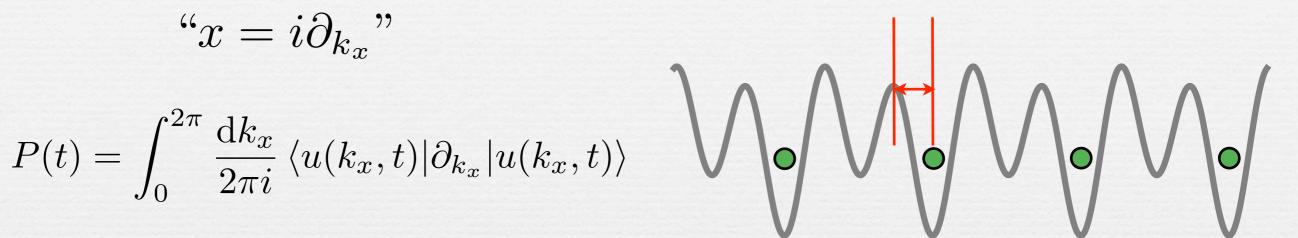


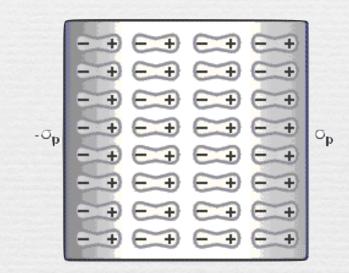
#### Polarization and Berry phase

Resta, King-Smith, Vanderbilt, ...

$$"x = i\partial_{k_x}"$$

$$P(t) = \int_0^{2\pi} \frac{\mathrm{d}k_x}{2\pi i} \langle u(k_x, t) | \partial_{k_x} | u(k_x, t) \rangle$$



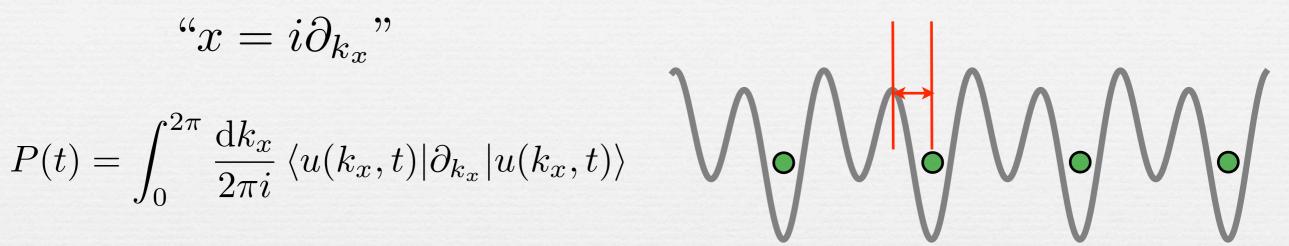


### Polarization and Berry phase

Resta, King-Smith, Vanderbilt, ...

$$"x = i\partial_{k_x}"$$

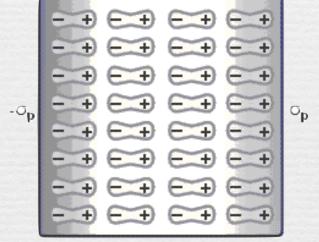
$$P(t) = \int_0^{2\pi} \frac{\mathrm{d}k_x}{2\pi i} \langle u(k_x, t) | \partial_{k_x} | u(k_x, t) \rangle$$



#### Change of Polarization

$$\Delta P = \int_0^T dP$$

$$= \frac{1}{2\pi i} \int_0^T \int_0^{2\pi} dt \, dk_x \left( \langle \partial_t u | \partial_{k_x} u \rangle - h.c. \right)$$
Berry Curvature



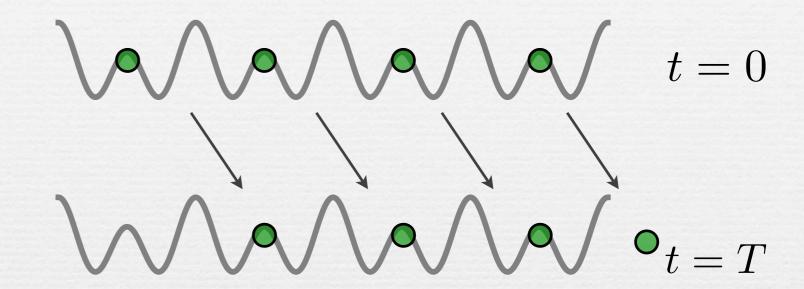
Berry Curvature

$$(\langle \partial_t u | \partial_{k_x} u \rangle - h.c.)$$

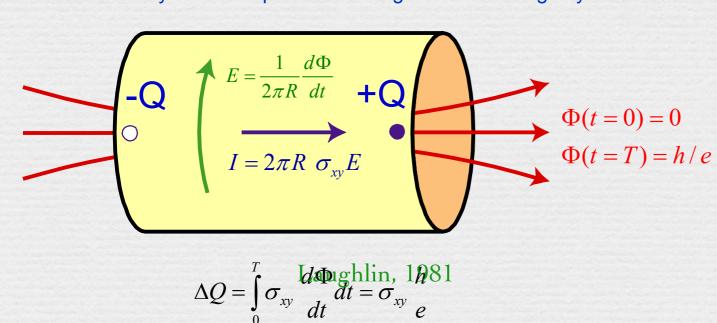
= Chern number

#### Connection to IQHE

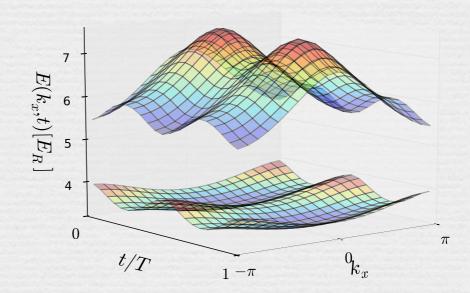
$$H(k_x, t) = H(k_x, t + T)$$



Adiabatically thread a quantum of magnetic flux through cylinder.

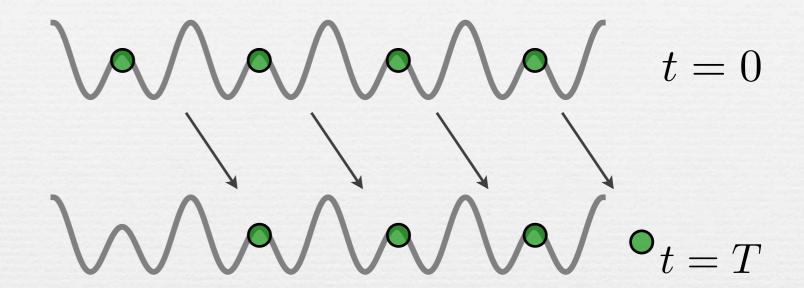


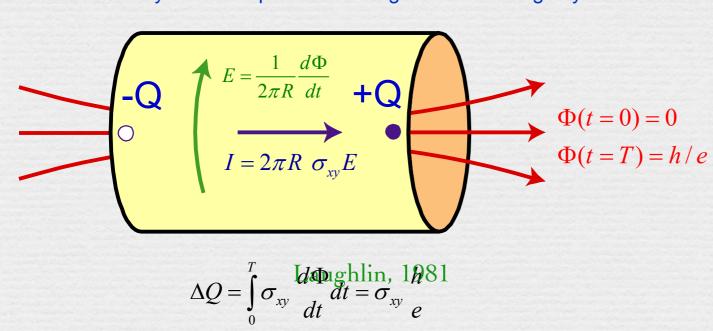
$$V_1 = 4E_R \quad V_2 = 4E_R$$



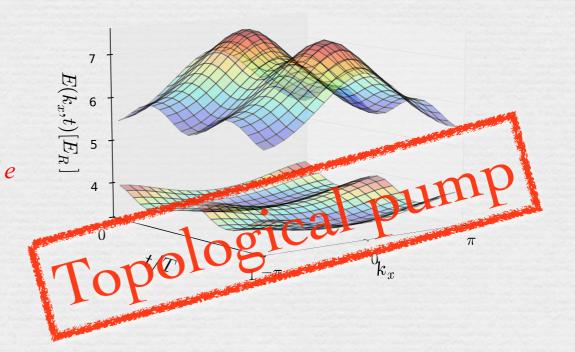
#### Connection to IQHE

$$H(k_x, t) = H(k_x, t + T)$$





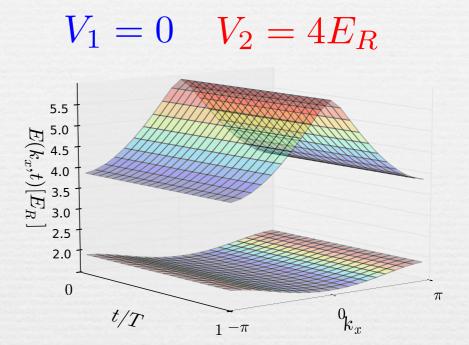
$$V_1 = 4E_R \quad V_2 = 4E_R$$



#### Adiabatic connection

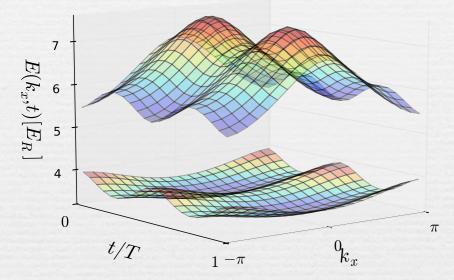
$$V_1 = 4E_R$$
  $V_2 = 4E_R$ 

$$V_1 \cos^2\left(\frac{2\pi x}{d}\right) + V_2 \cos^2\left(\frac{\pi x}{d} - \frac{\pi t}{T}\right)$$



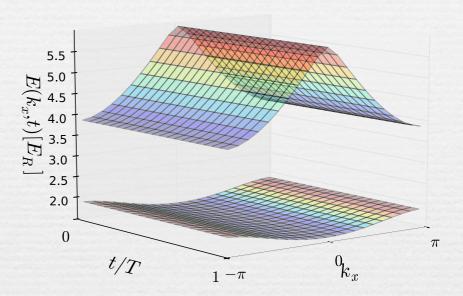
#### Adiabatic connection

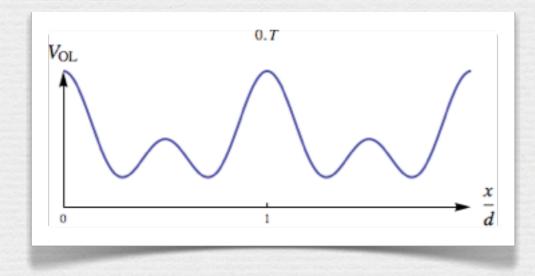
$$V_1 = 4E_R \quad V_2 = 4E_R$$



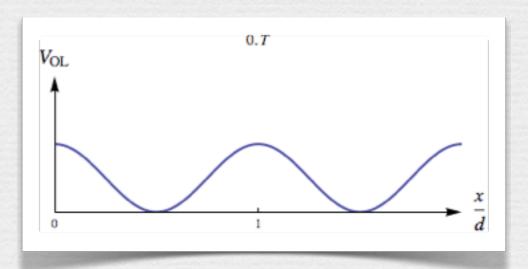
$$V_1 \cos^2\left(\frac{2\pi x}{d}\right) + V_2 \cos^2\left(\frac{\pi x}{d} - \frac{\pi t}{T}\right)$$

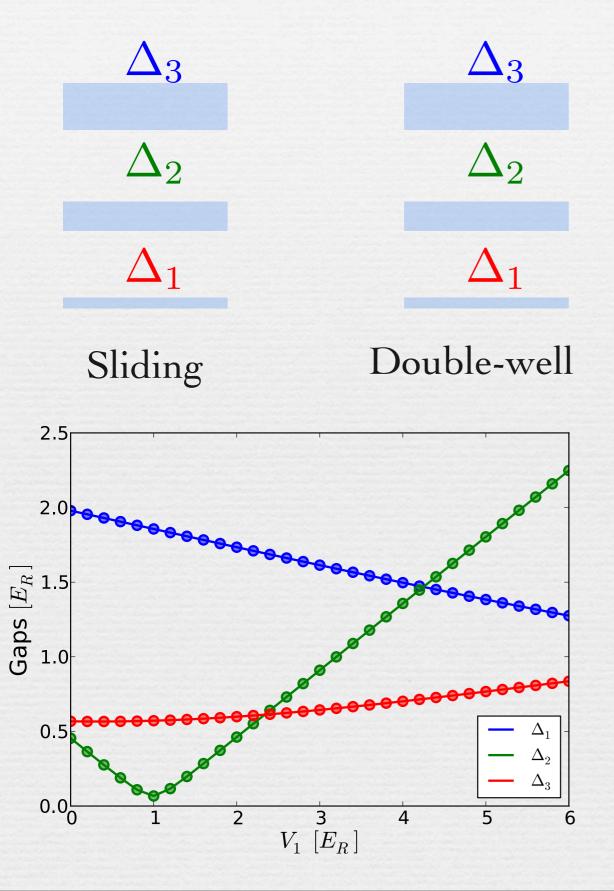


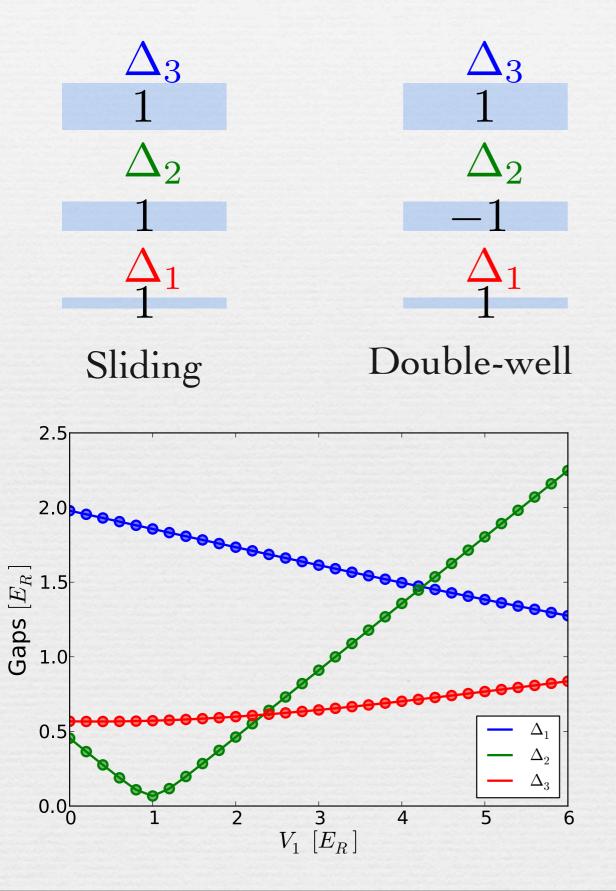


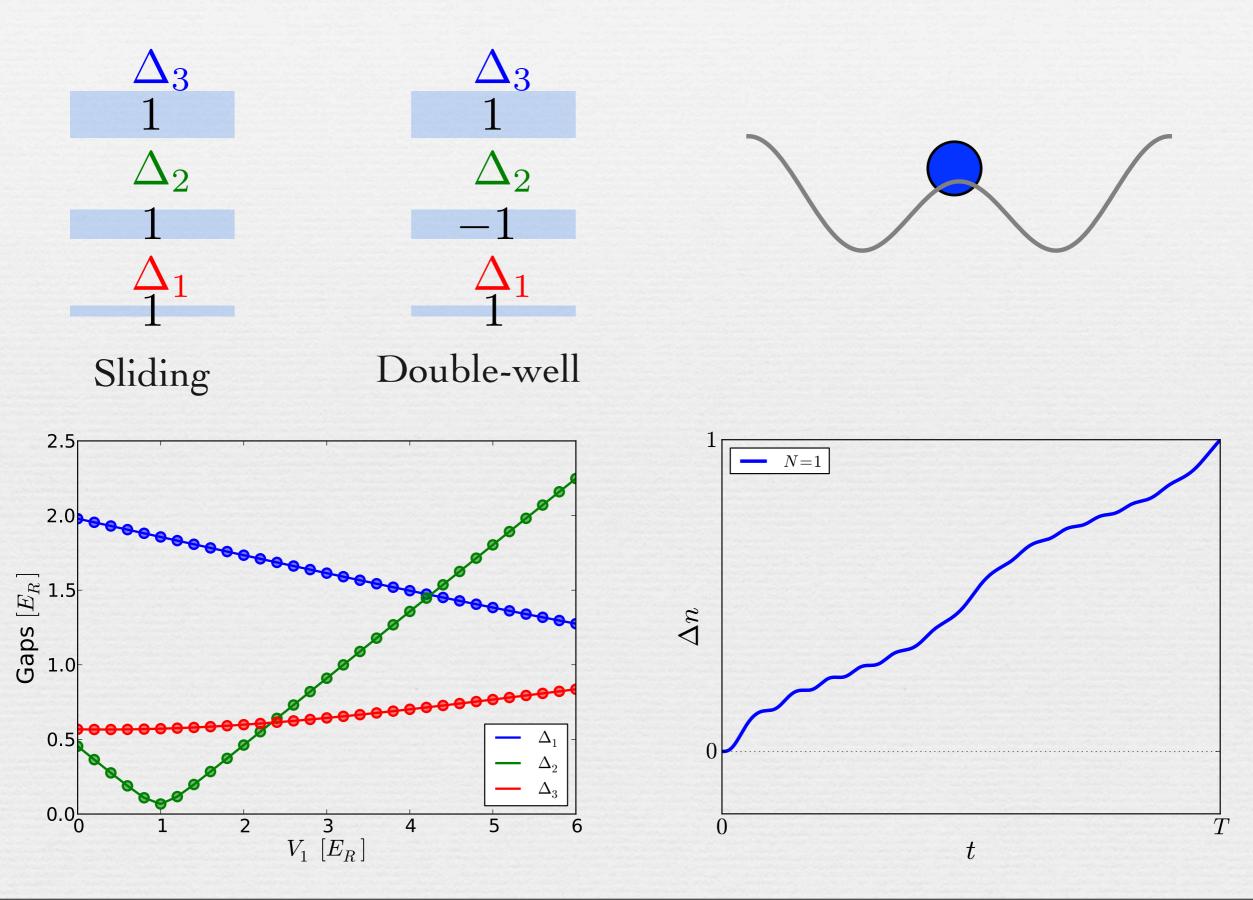


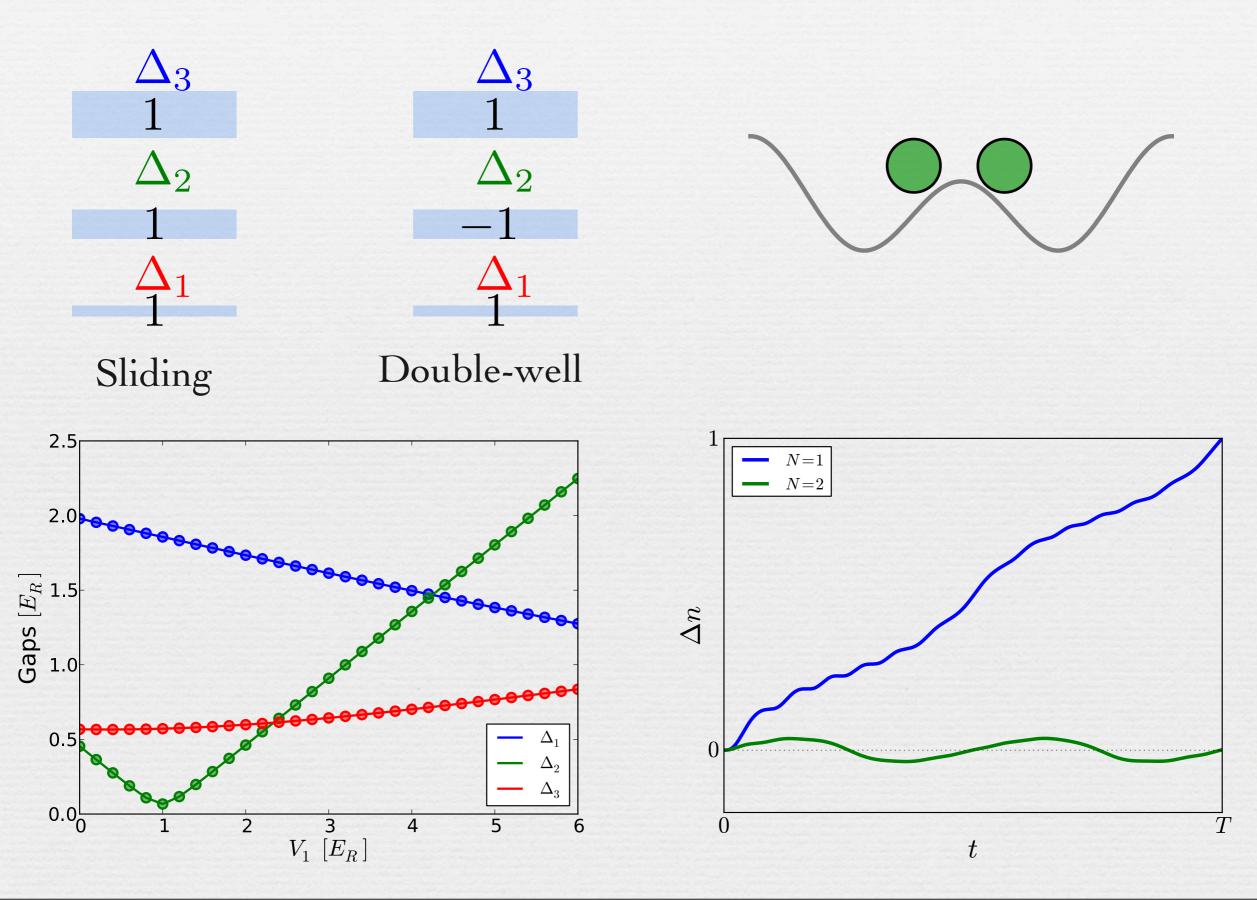


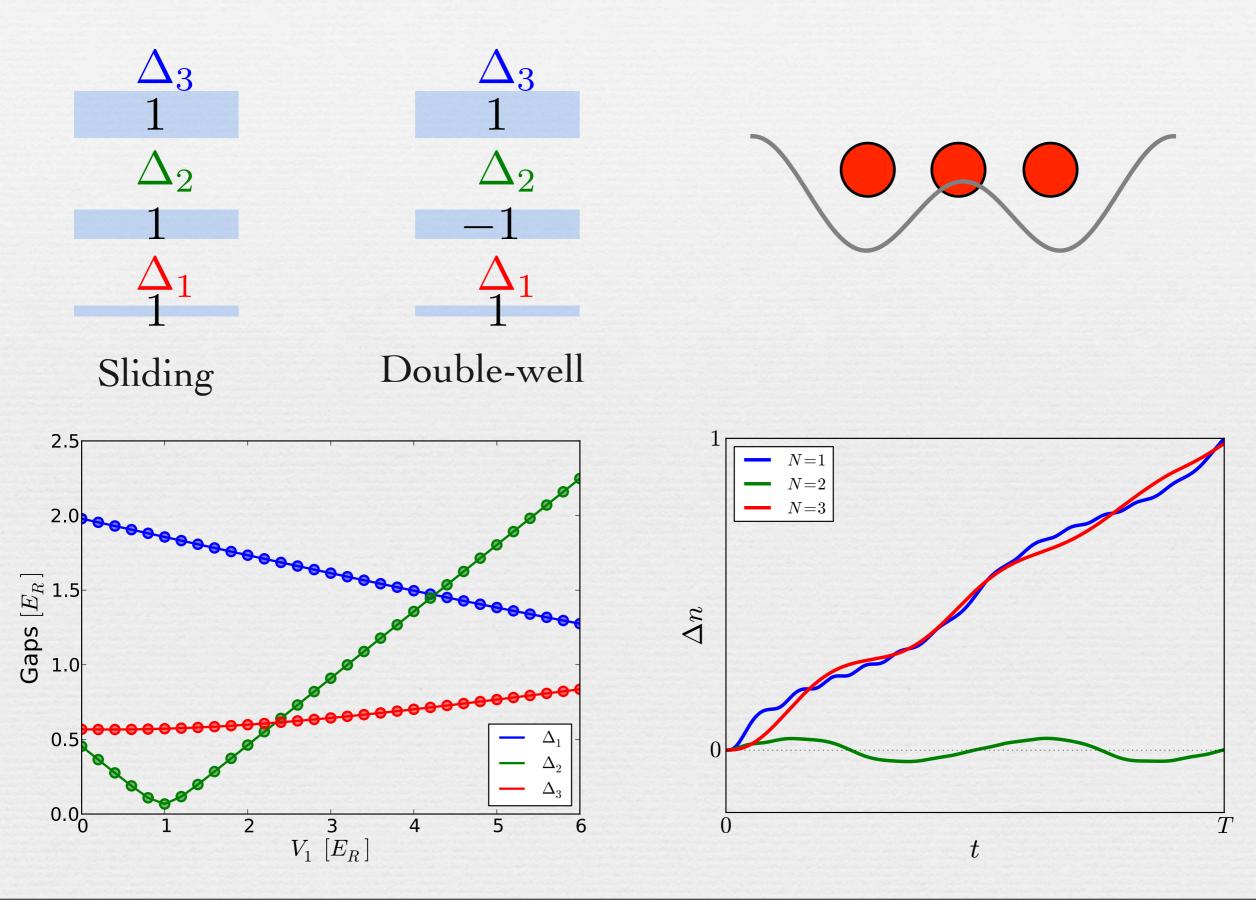




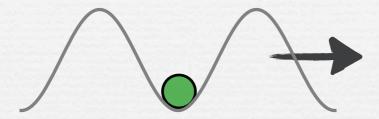








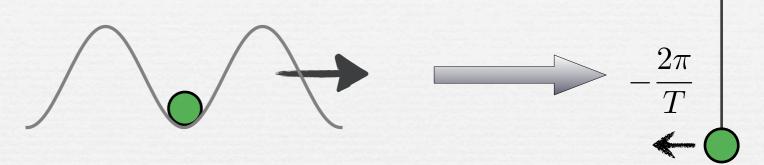
$$\left(m\ddot{x} = -\frac{\partial V_{\mathrm{OL}}(x,t)}{\partial x}\right)$$



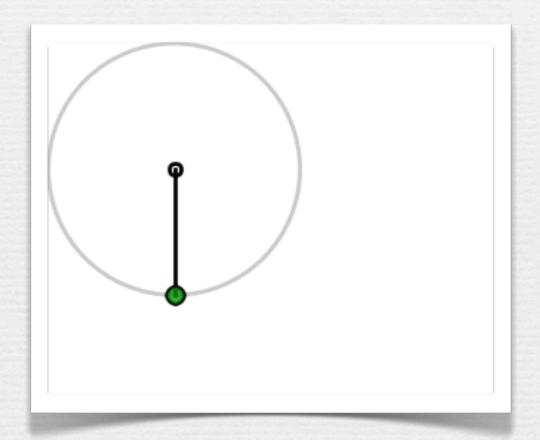
$$m\ddot{x} = -\frac{\partial V_{\rm OL}(x,t)}{\partial x} \qquad -\frac{2\pi}{T}$$

Without V<sub>1</sub> term, pumping maps to a simple pendulum

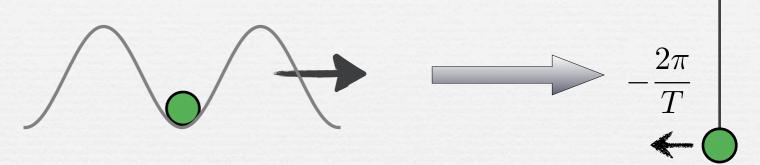
$$m\ddot{x} = -\frac{\partial V_{\mathrm{OL}}(x,t)}{\partial x}$$



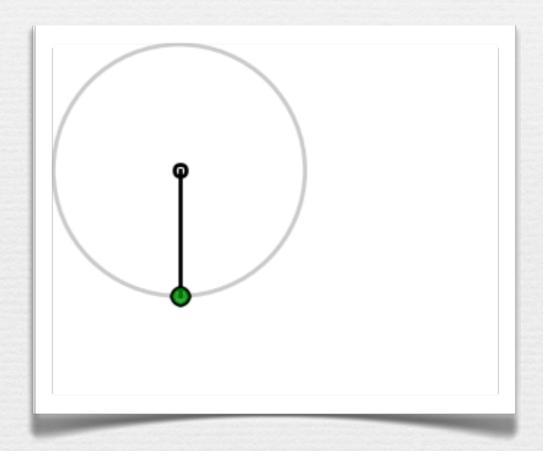
- Without V<sub>1</sub> term, pumping maps to a simple pendulum
  - Slow pumping-> small oscillations-> follows pumping



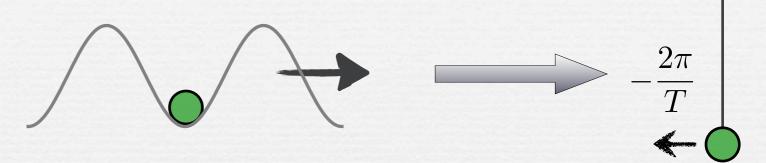
$$m\ddot{x} = -\frac{\partial V_{\rm OL}(x,t)}{\partial x}$$



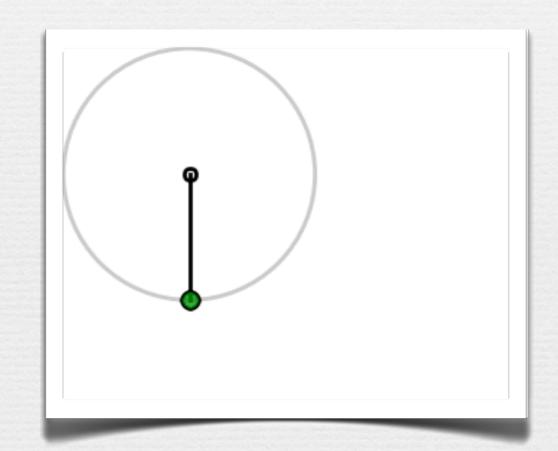
- Without V<sub>1</sub> term, pumping maps to a simple pendulum
  - Slow pumping-> small oscillations-> follows pumping
  - Fast pumping-> swings around pivot-> can not follow pumping



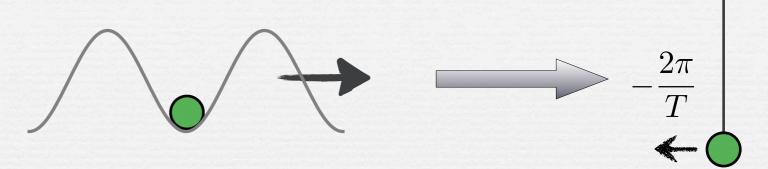
$$m\ddot{x} = -\frac{\partial V_{\rm OL}(x,t)}{\partial x}$$



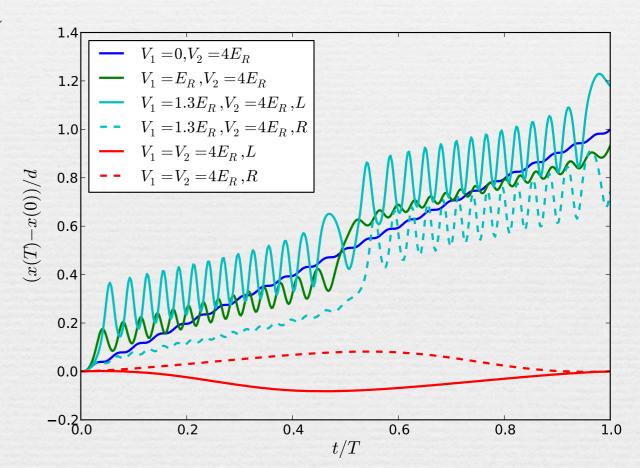
- Without V<sub>1</sub> term, pumping maps to a simple pendulum
  - Slow pumping-> small oscillations-> follows pumping
  - Fast pumping-> swings around pivot-> can not follow pumping



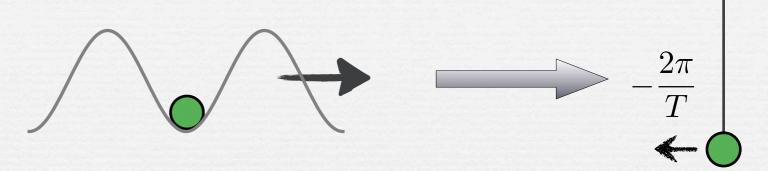
$$m\ddot{x} = -\frac{\partial V_{\mathrm{OL}}(x,t)}{\partial x}$$



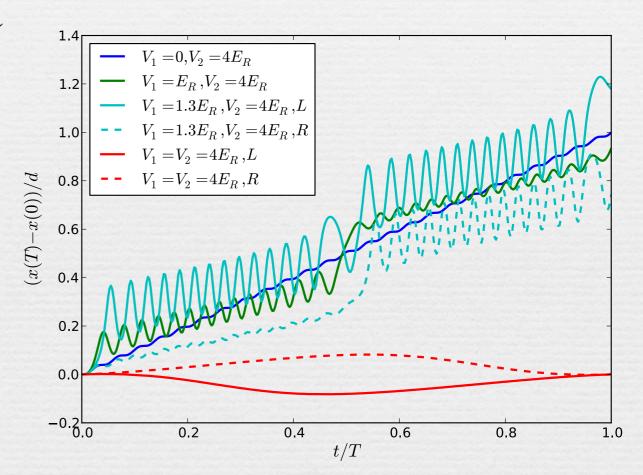
- Without V<sub>1</sub> term, pumping maps to a simple pendulum
  - Slow pumping-> small oscillations-> follows pumping
  - Fast pumping-> swings around pivot-> can not follow pumping
- V<sub>1</sub> term: Driven pendulum shows chaotic behavior.



$$m\ddot{x} = -\frac{\partial V_{\mathrm{OL}}(x,t)}{\partial x}$$



- Without V<sub>1</sub> term, pumping maps to a simple pendulum
  - Slow pumping-> small oscillations-> follows pumping
  - Fast pumping-> swings around pivot-> can not follow pumping
- V<sub>1</sub> term: Driven pendulum shows chaotic behavior.



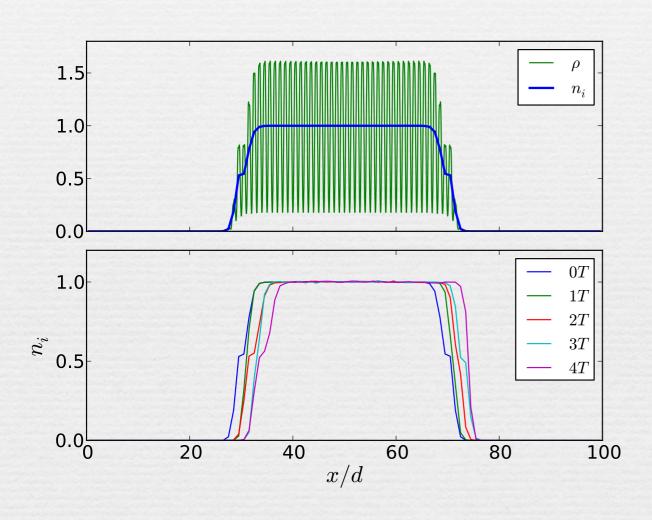
In general, classical pumped charge is not quantized

## Practical issues

- Detection
- External trap
- Temperature effect
- Non-adiabatic effect

## Trapping & Detection

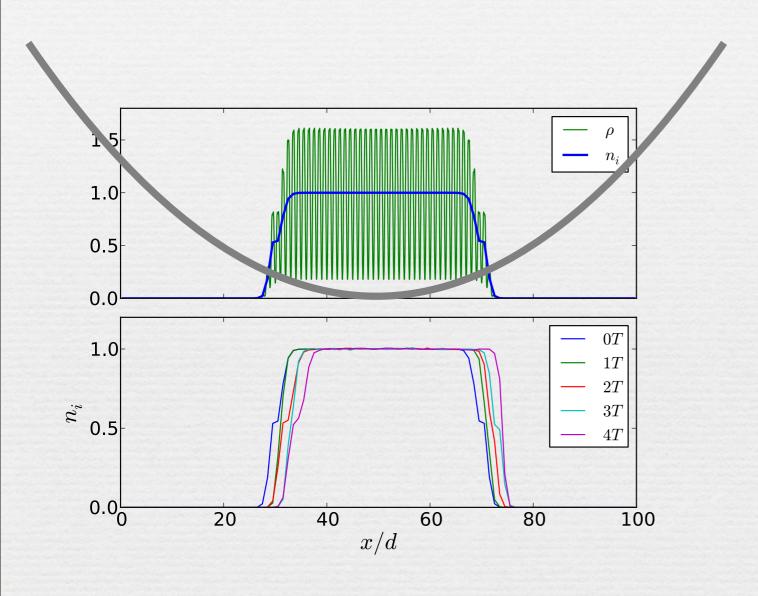
LW, Troyer and Dai, 1301.7435



$$\langle x \rangle / d = \Delta n$$

## Trapping & Detection

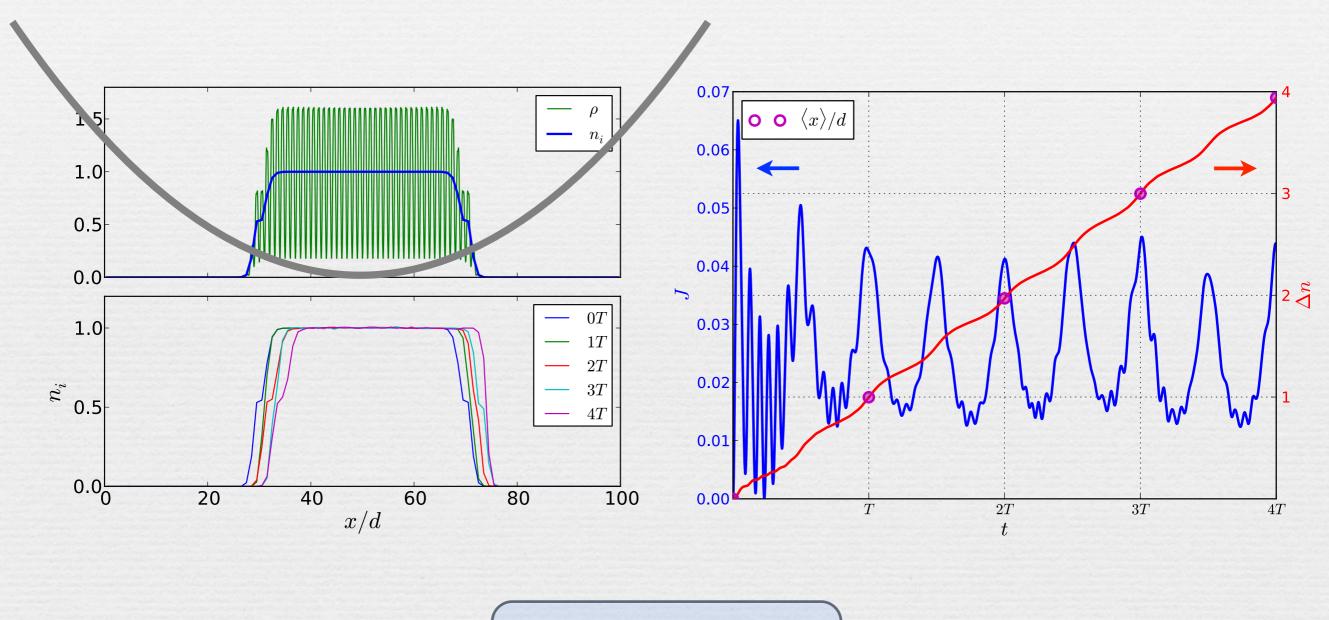
LW, Troyer and Dai, 1301.7435



$$\langle x \rangle / d = \Delta n$$

### Trapping & Detection

LW, Troyer and Dai, 1301.7435



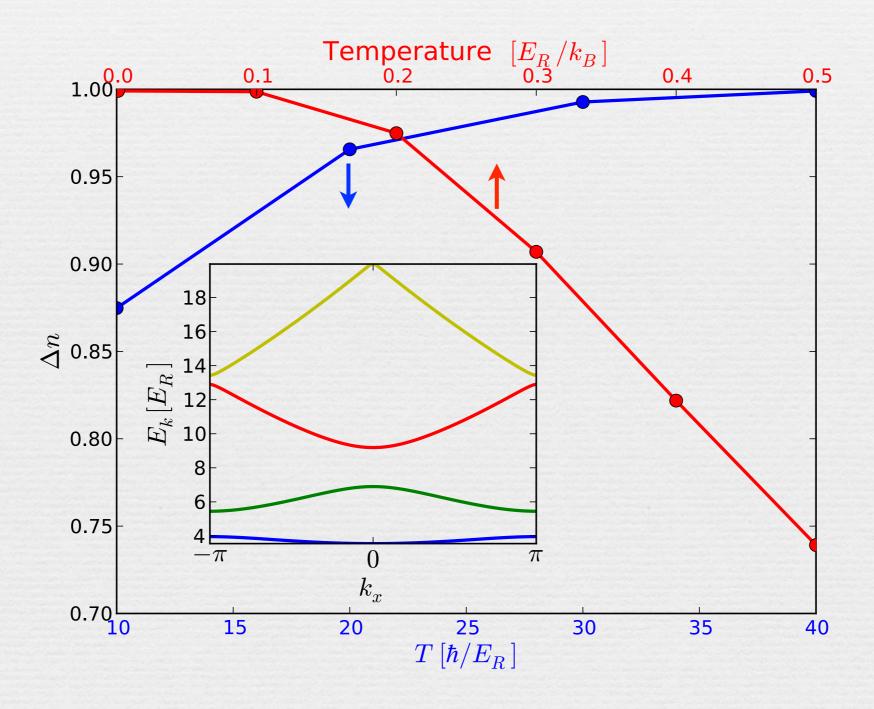
$$\langle x \rangle / d = \Delta n$$

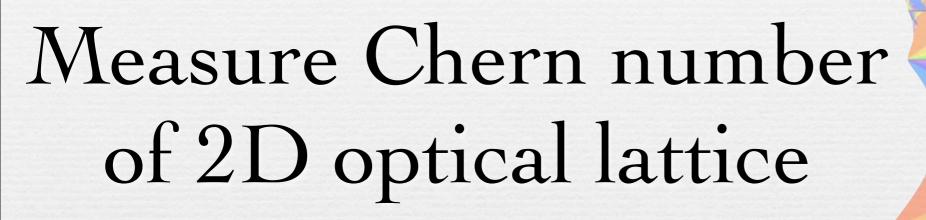
#### Temperature & Non-adiabatic effect

Temperature 
$$\ll \frac{\Delta}{k_B}$$
  $T \gg \frac{\hbar}{\Delta}$ 

#### Temperature & Non-adiabatic effect

Temperature 
$$\ll \frac{\Delta}{k_B}$$
  $T \gg \frac{\hbar}{\Delta}$ 



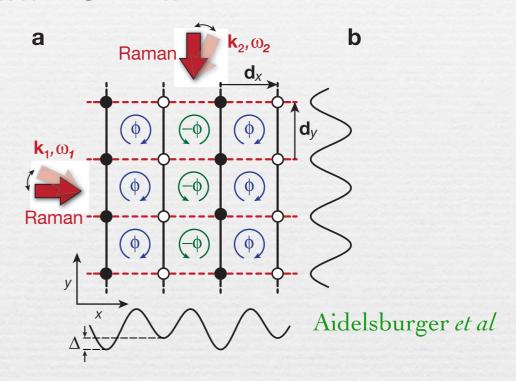


with

# Synthetic gauge-field in optical lattices

Imprint complex phases to the hopping amplitude

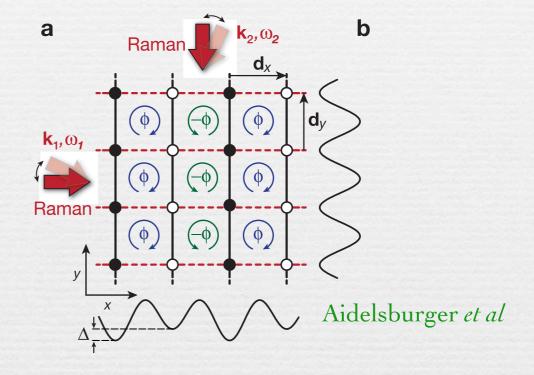
Staggered flux lattice Munich



# Synthetic gauge-field in optical lattices

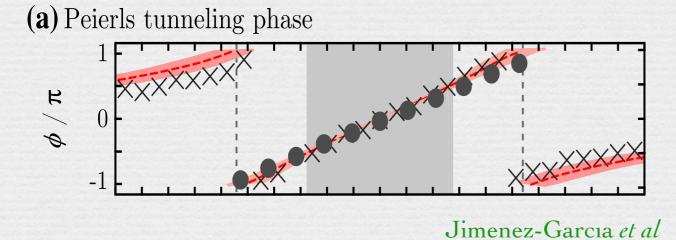
Imprint complex phases to the hopping amplitude

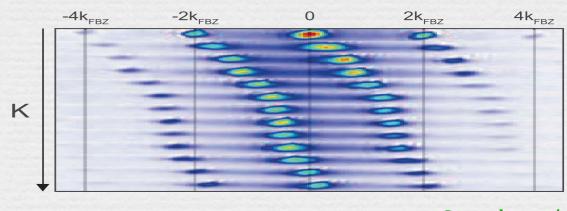
Staggered flux lattice Munich



1D Peierls lattice NIST, Hamburg

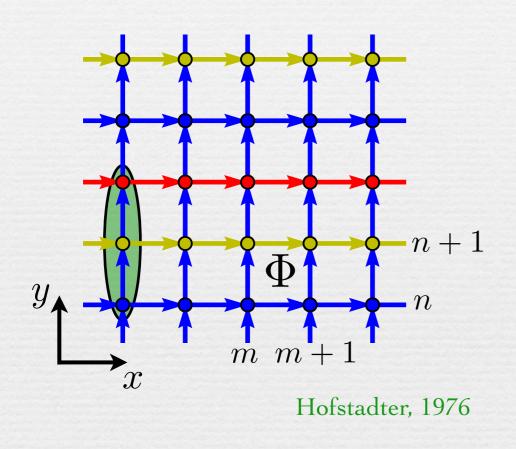
$$H = -J \sum_{m} e^{i2\pi\Phi} c_{m+1}^{\dagger} c_m + H.c.$$

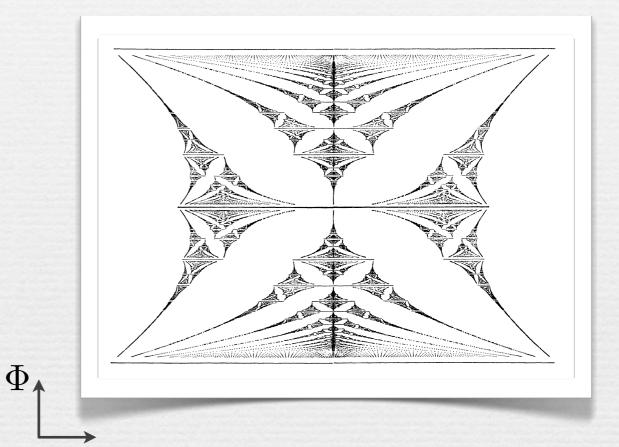




## Hofstadter optical lattice

$$H = -J \sum_{m,n} e^{i2\pi n\Phi} c_{m+1,n}^{\dagger} c_{m,n} + c_{m,n+1}^{\dagger} c_{m,n} + H.c. \quad \Phi = p/q$$

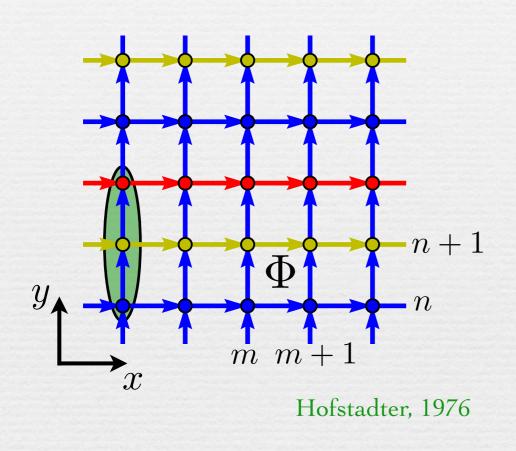


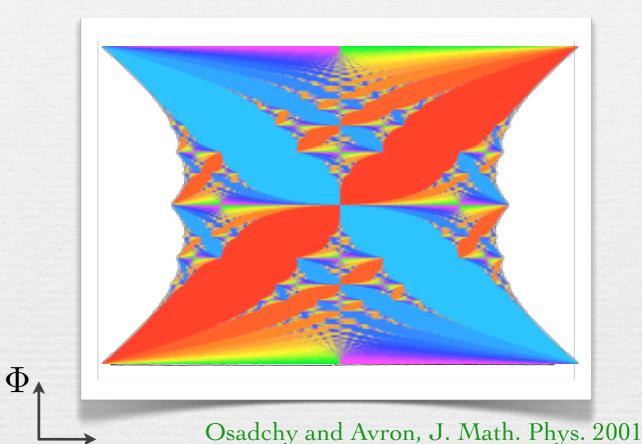




## Hofstadter optical lattice

$$H = -J \sum_{m,n} e^{i2\pi n\Phi} c_{m+1,n}^{\dagger} c_{m,n} + c_{m,n+1}^{\dagger} c_{m,n} + H.c. \quad \Phi = p/q$$



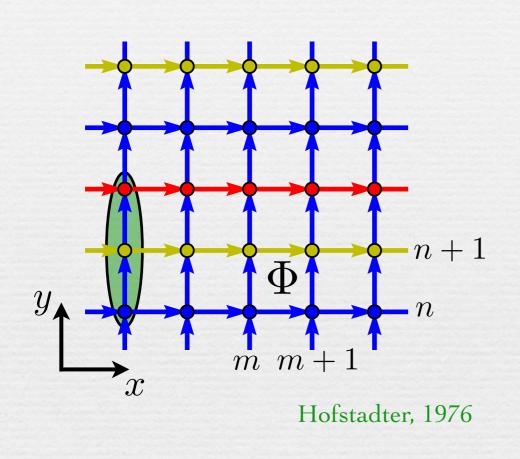


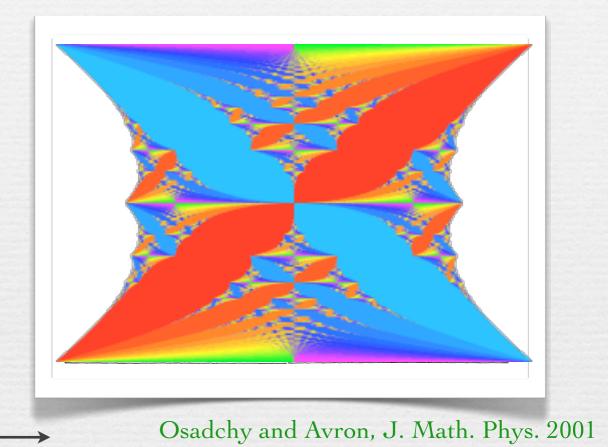


## Hofstadter optical lattice

$$H = -J \sum_{m,n} e^{i2\pi n\Phi} c_{m+1,n}^{\dagger} c_{m,n} + c_{m,n+1}^{\dagger} c_{m,n} + H.c. \quad \Phi = p/q$$

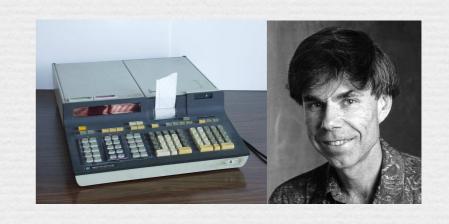
 $\Phi_{\bullet}$ 





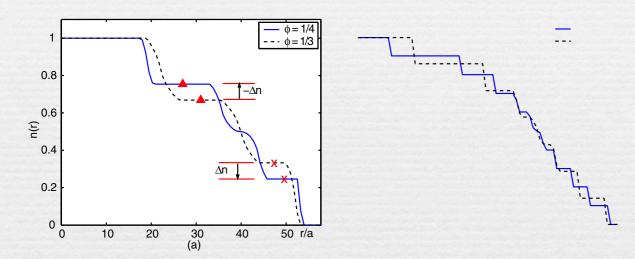
NO sharp edge states in harmonic trapping potential

Buchhold et al



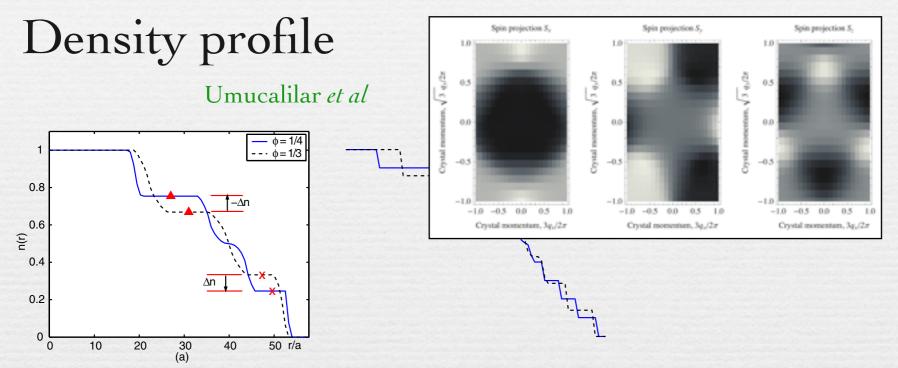
#### Density profile

Umucalilar et al



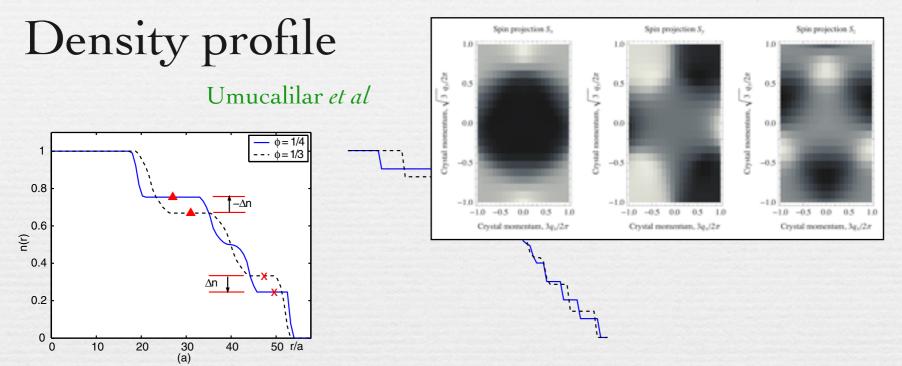
Time-of-flight

Alba et al, Zhao et al



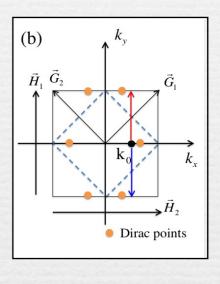
Time-of-flight

Alba et al, Zhao et al



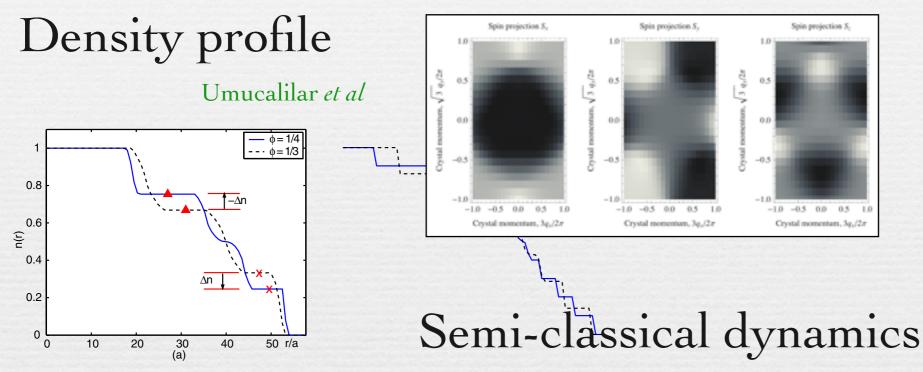
#### Zak phases

Abanin et al



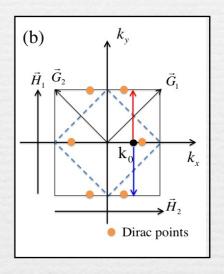
#### Time-of-flight

Alba et al, Zhao et al



Zak phases

Abanin et al

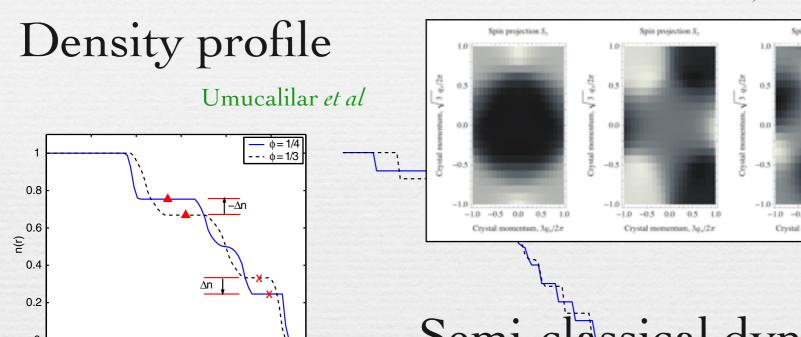


Price et al

$$\mathbf{\dot{r}} = \frac{1}{\hbar} \frac{\partial \varepsilon(\mathbf{k})}{\partial \mathbf{k}} - \frac{\mathbf{F}}{\hbar} \times \hat{\mathbf{z}} \Omega(\mathbf{k})$$

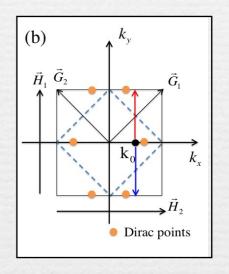
Time-of-flight

Alba et al, Zhao et al



Zak phases

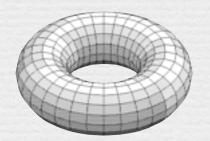
Abanin et al



Semi-classical dynamics

Price et al

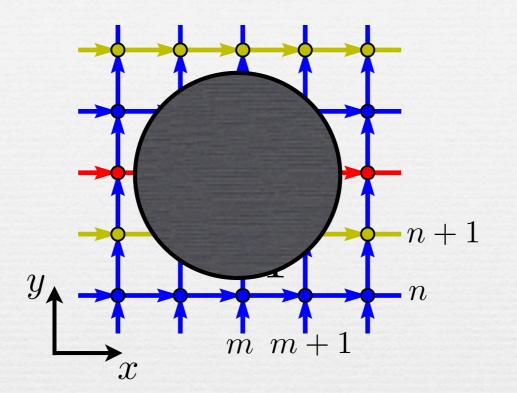
$$\mathbf{\dot{r}} = \frac{1}{\hbar} \frac{\partial \varepsilon(\mathbf{k})}{\partial \mathbf{k}} - \frac{\mathbf{F}}{\hbar} \times \hat{\mathbf{z}} \Omega(\mathbf{k})$$

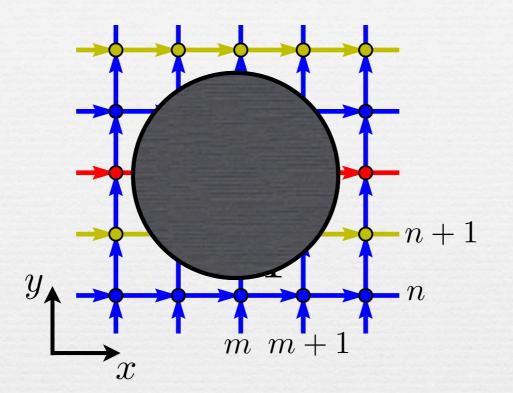


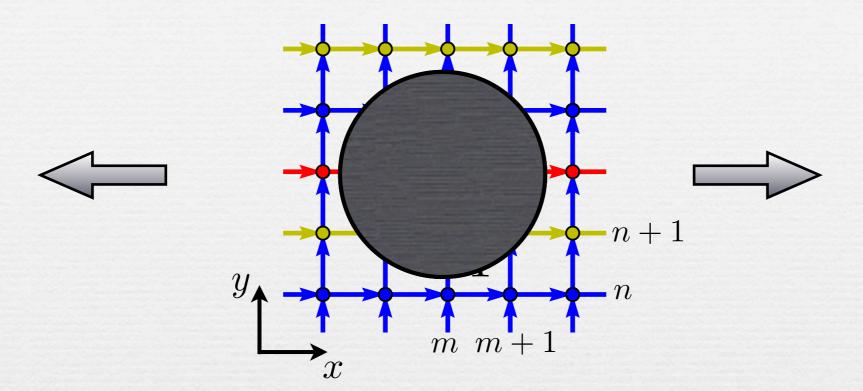
We propose a new probe based on

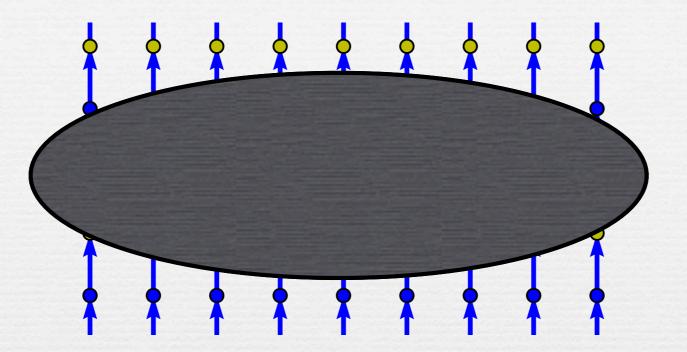
Topological Pumping Effect

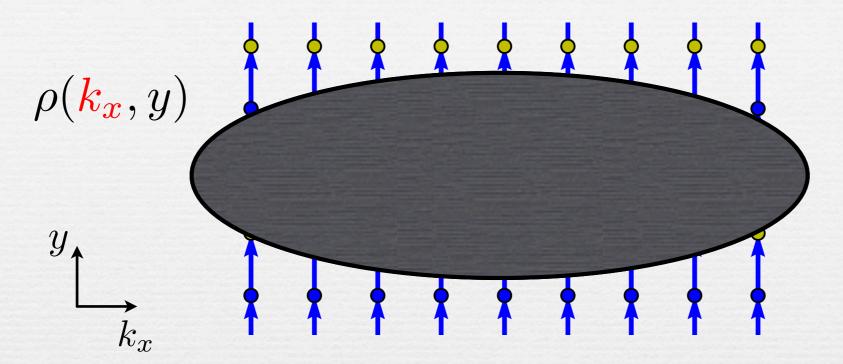
 $ho(\pmb{k_x},y)$ 

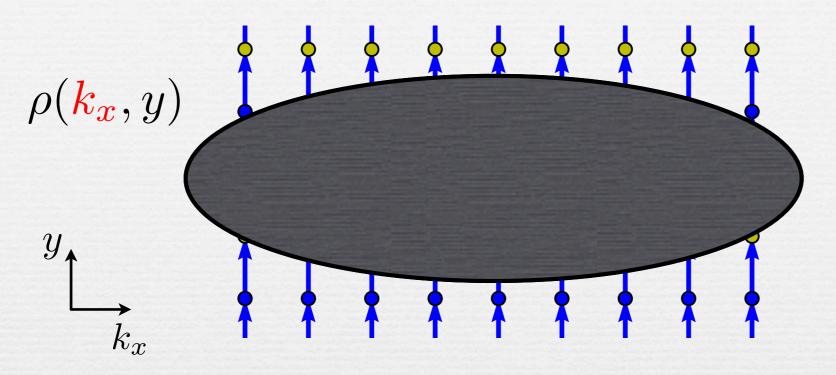




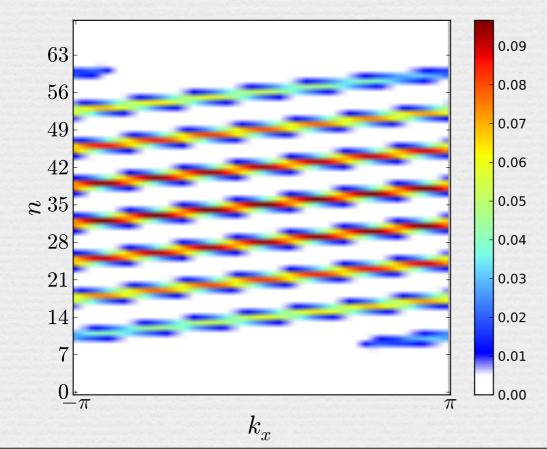


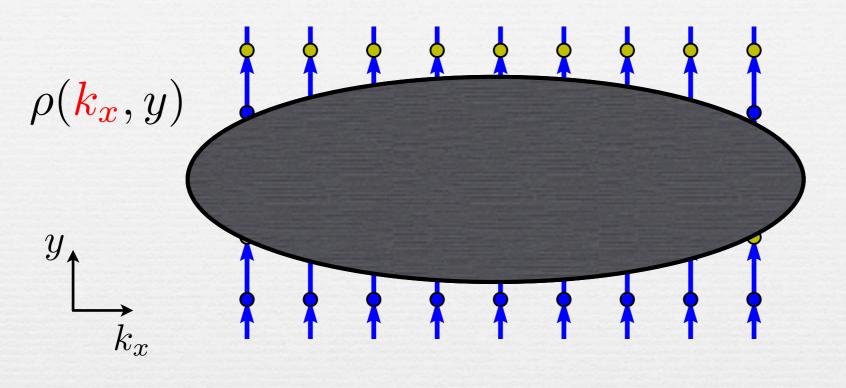


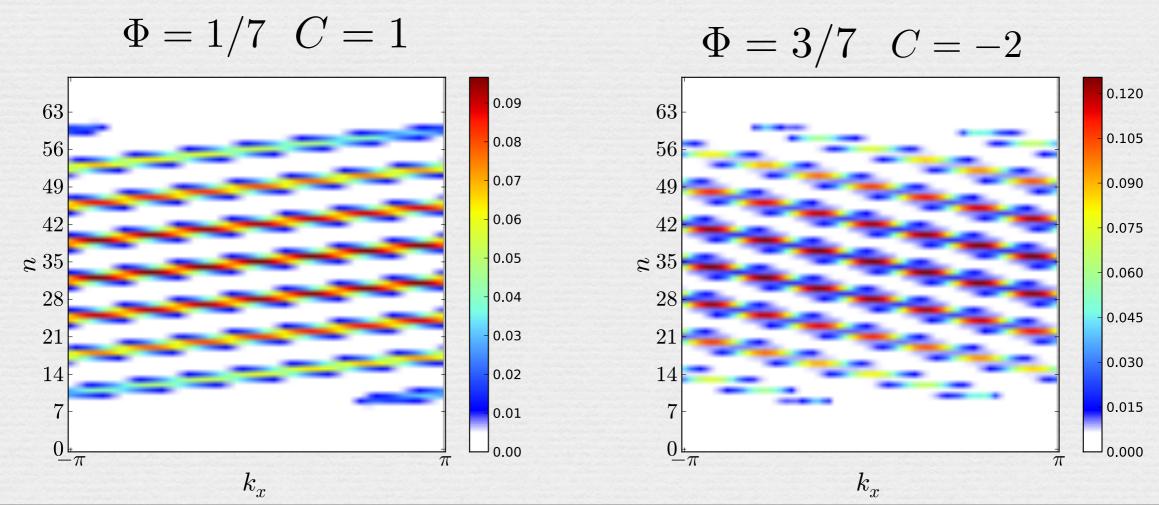


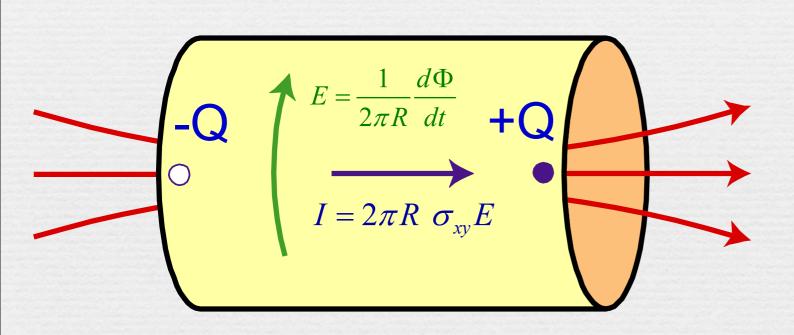


$$\Phi = 1/7 \ C = 1$$



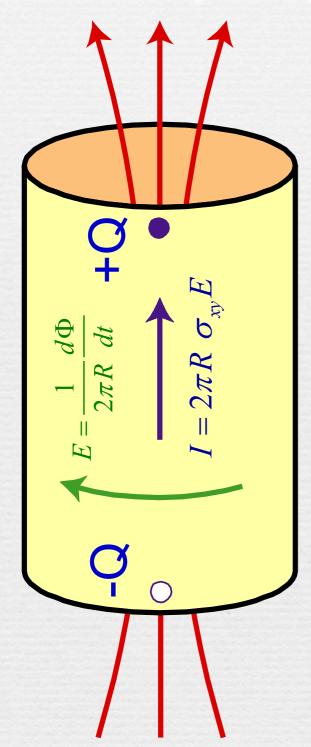


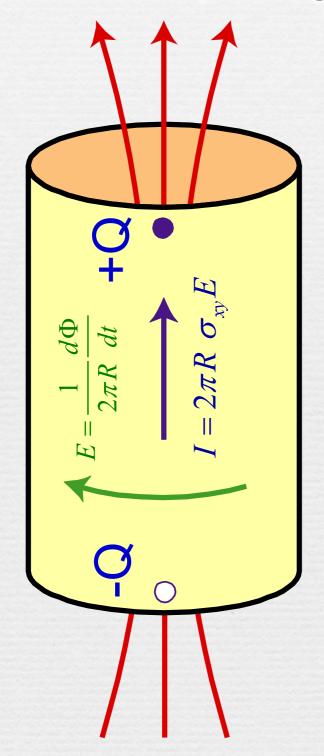


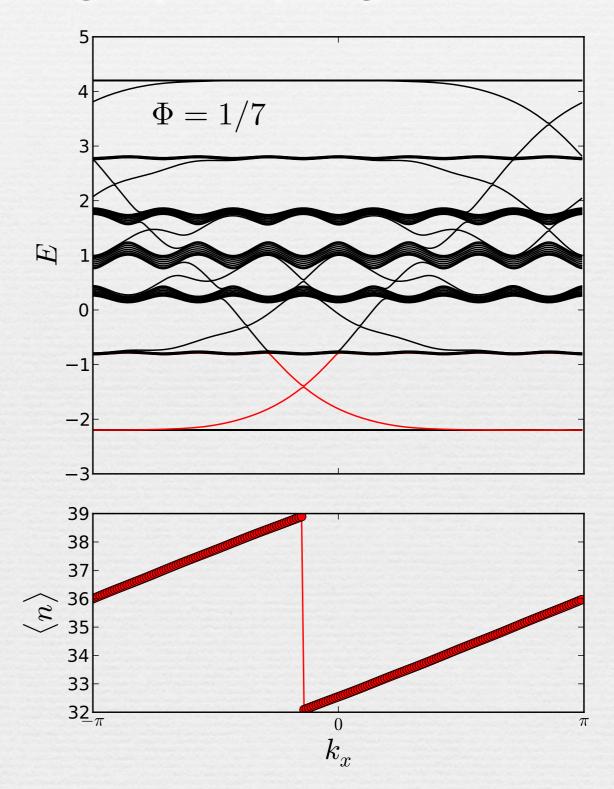


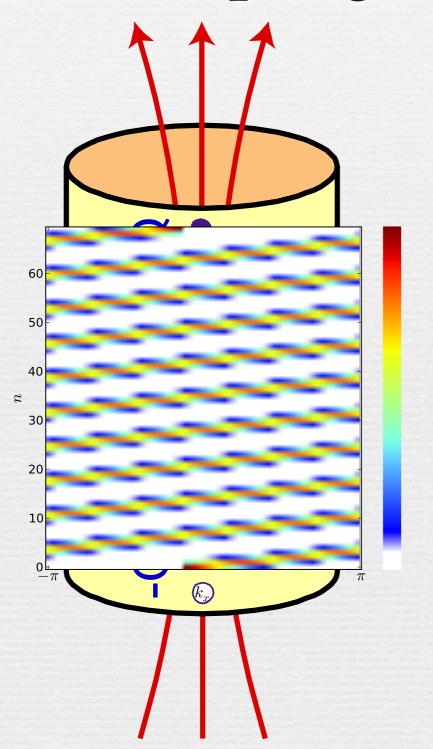
$$\Delta Q = \int_{0}^{T} \sigma_{xy} \frac{d\Phi}{dt} dt = \sigma_{xy} \frac{h}{e}$$

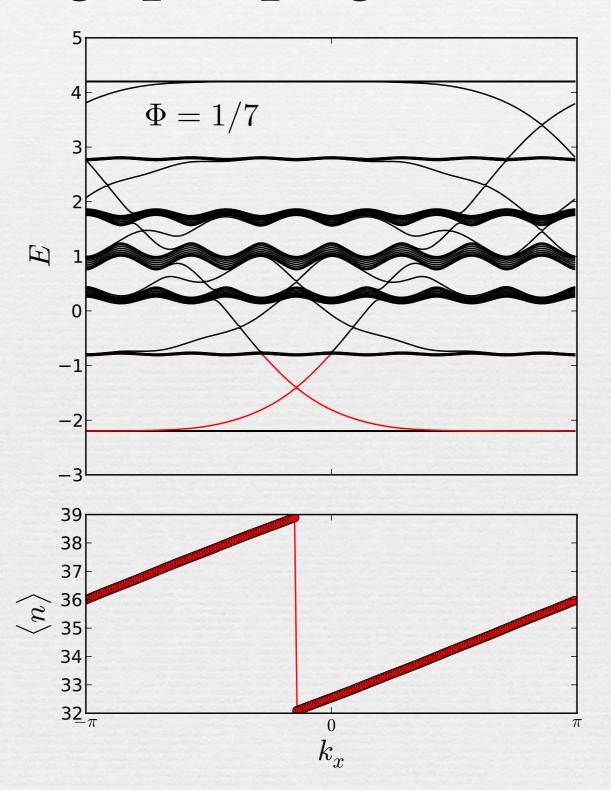
$$\Delta Q = ne \rightarrow \sigma_{xy} = n \frac{e^2}{h}$$





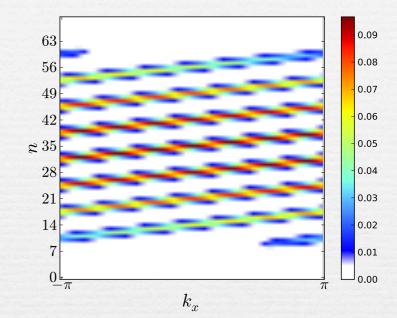






#### Quantitative Characterizations

- Slope
- \* # of cuts (edge modes)
- COM along y-direction



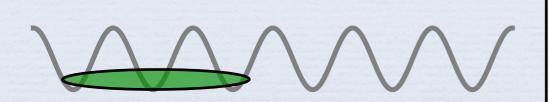
Bi-partition number of particle (trace index)

#### Salient features

- Bulk detection, does not require edge states
- $\rho(k_x, y)$  is nearly impossible to measure in solids, but accessible to cold atom toolbox
- Can be extended to interacting case

#### Fractional charge pumping

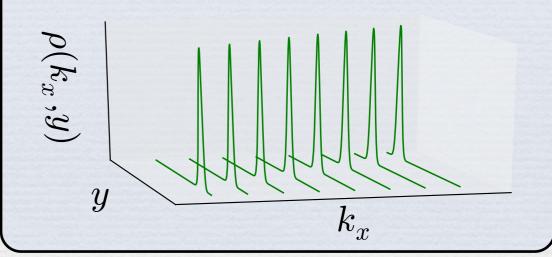
1D lattice



Interaction is crucial for opening an energy gap

2D Laughlin state

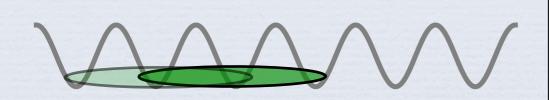
$$\rho(k_x, y) = \frac{\nu}{\sqrt{\pi}} e^{-(y - k_x)^2}$$



Can be used to detect FQHE and fractional Chern insulators realized in optical lattice Cooper et al, Yao et al, Nielsen et al

### Fractional charge pumping

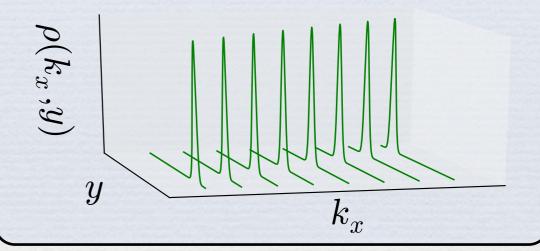
1D lattice



Interaction is crucial for opening an energy gap

2D Laughlin state

$$\rho(k_x, y) = \frac{\nu}{\sqrt{\pi}} e^{-(y - k_x)^2}$$



Can be used to detect FQHE and fractional Chern insulators realized in optical lattice Cooper et al, Yao et al, Nielsen et al

#### Laughlin states on lattice

 $\Phi_{\text{Laughlin}} = \Phi_{\text{Hofstadter}}^{1/\nu}$ 

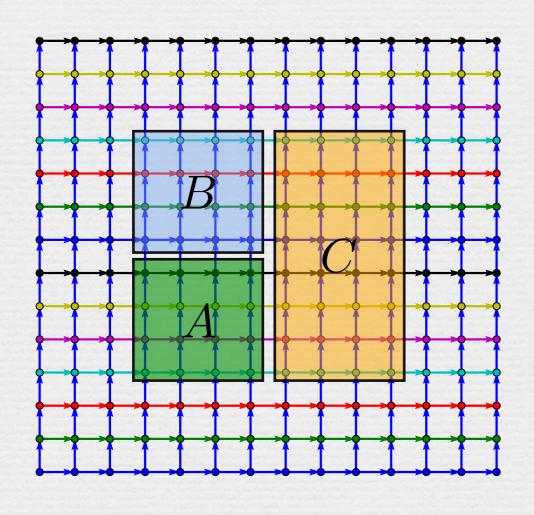
#### Laughlin states on lattice

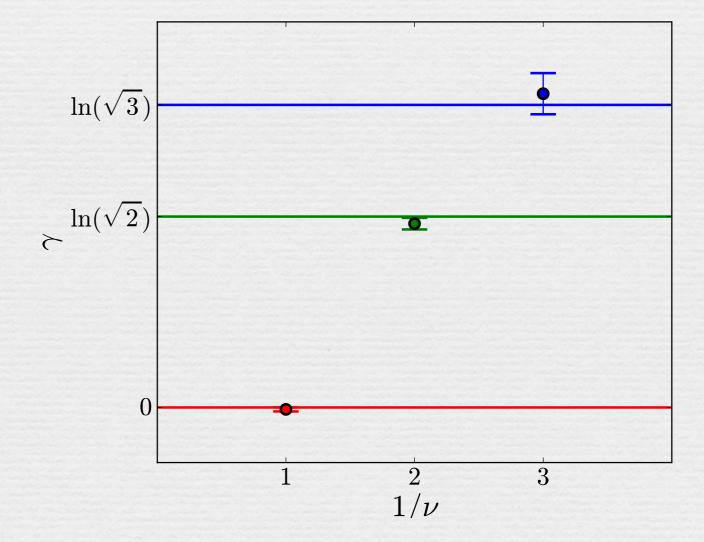
$$\Phi_{\text{Laughlin}} = \Phi_{\text{Hofstadter}}^{1/\nu}$$

Topological entanglement entropy

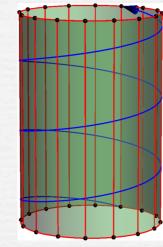
Preskill, Kitaev, Levin, Wen Zhang *et al* 

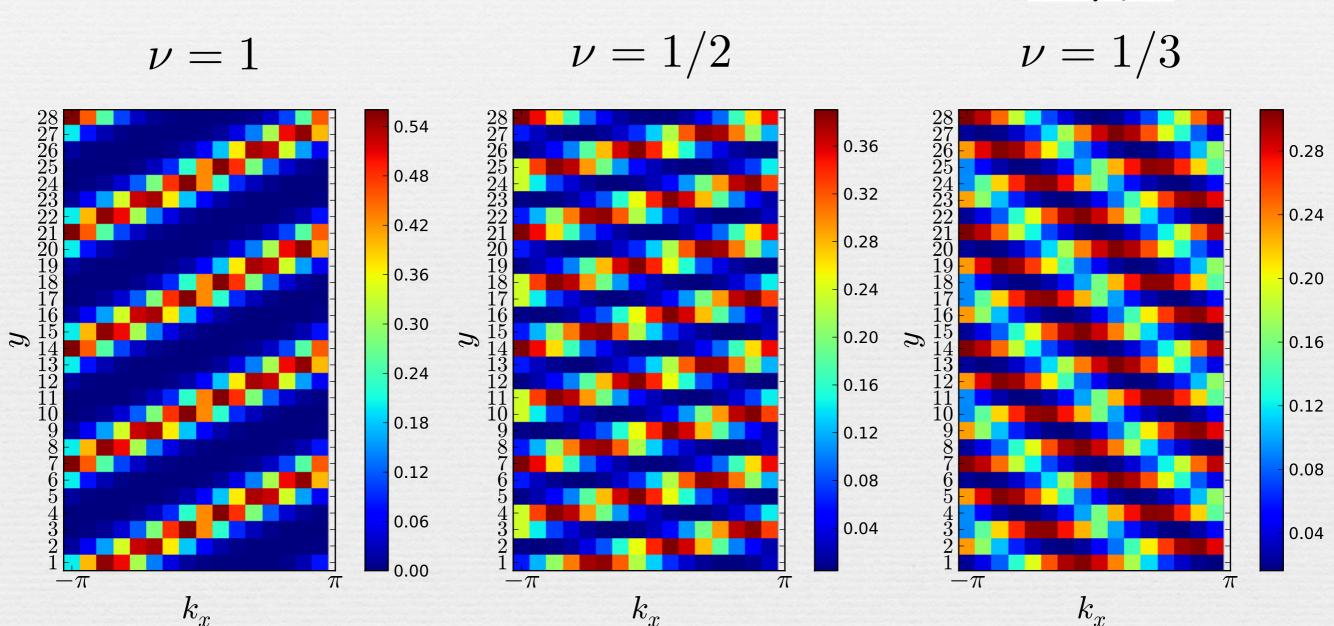
$$-\gamma = S_A + S_B + S_C - S_{AB} - S_{BC} - S_{AC} + S_{ABC}$$





### Hybrid densities

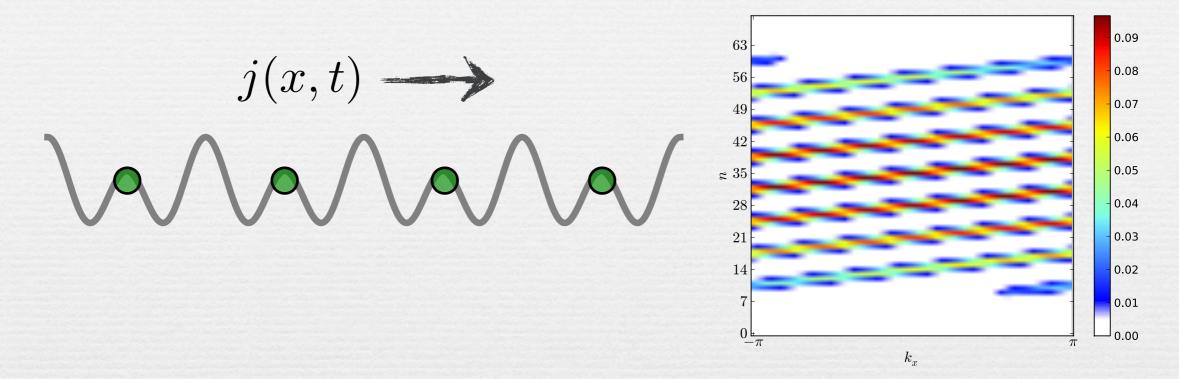




HTOF is also useful to detect FQHE state!

#### Summary

arXiv:1301.7435 PRL 110, 166802

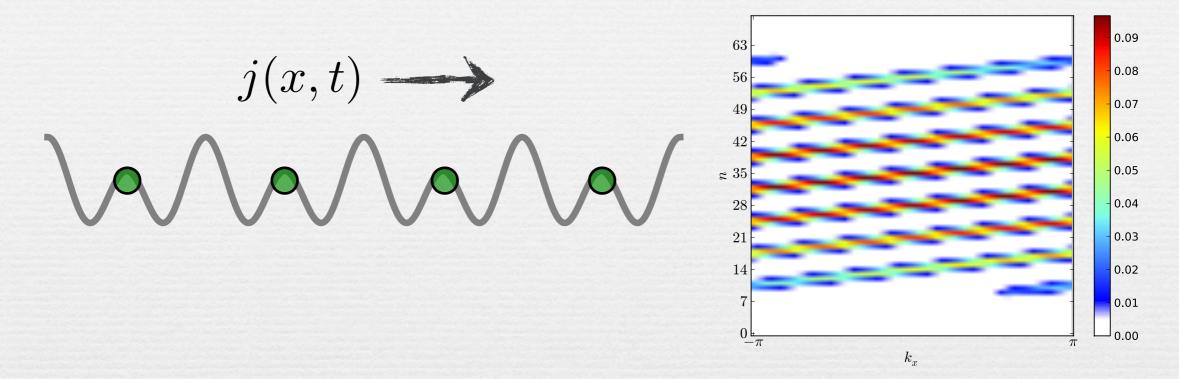


Topological charge pumping is a common thread unifies many features of topological states

Guideline for design and detection of topological phases in cold atom systems

#### Summary

arXiv:1301.7435 PRL 110, 166802



Topological charge pumping is a common thread unifies many features of topological states

Guideline for design and detection of topological phases in cold atom systems

#### You might try it in your lab!

# Thank you!